

Lecture I . Heat equation :

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Let $u(t,x)$ denote the temperature at point x at time t .

One dimensional heat equation is:

$$\partial_t u = k \partial_x^2 u$$

where $k > 0$ is a constant (the thermal conductivity of the material)

That is, the change in heat at a specific point is proportional to the second derivative of the heat along the wire.

• Method of separation of variables

we write $u(t,x) = z(t) y(x)$, a product of a function of t and a function of x .

we obtain, $\partial_t u(t,x) = z'(t) y(x)$

$$\partial_x^2 u(t,x) = z(t) y''(x)$$

Thus, $z'(t) y(x) = k z(t) y''(x)$

assuming $z(t), y(x)$ are non-zero,

we have $\frac{z'(t)}{z(t)} = \frac{y''(x)}{y(x)} = \lambda \quad k=1$

Since the left-hand ^{side} is a constant with respect to x ,

and the right-hand side is a constant w.r.t t ,

both sides must be constant.

$$\Rightarrow \begin{cases} z'(t) - \lambda z(t) = 0 \\ y''(x) - \lambda y(x) = 0 \end{cases}$$

$\lambda < 0$: physical solution

$$\text{Let } \lambda = -\omega^2. \text{ Then, } \begin{cases} z'(t) + \omega^2 z(t) = 0 & \textcircled{1} \\ y''(x) + \omega^2 y(x) = 0 & \textcircled{2} \end{cases}$$

The general solution of $\textcircled{2}$ takes the form

$$y(x) = a_\omega \cos(\omega x) + b_\omega \sin(\omega x)$$

The general solution of $\textcircled{1}$ is

$$z(t) = c \exp(-\omega^2 t)$$

So we have $u(t, x) = z(t) y(x)$

$$= \exp(-\omega^2 t) (A_\omega \cos(\omega x) + B_\omega \sin(\omega x))$$

observe that superposition (linear combination) is also a solution. Thus.

$$\begin{aligned} u(t, x) &= c_0 + \sum_{\omega=0}^{\infty} \exp(-\omega^2 t) (A_\omega \cos(\omega x) + B_\omega \sin(\omega x)) \\ &= c_0 + \sum_{\omega=0}^{\infty} (a_\omega \cos(\omega x) + b_\omega \sin(\omega x)) \end{aligned}$$

solve the equation.

Note that it is a representation in the form of a Fourier series with coefficients depending on the time t

II. wave equation

method of separation of variables:

$$\text{(eq)} \quad \partial_t^2 \Psi(t, x) = \partial_x^2 \Psi(t, x)$$

$$\text{Let } \Psi(t, x) = \underset{\substack{\uparrow \\ \text{shape}}}{A(x)} \cdot \underset{\substack{\uparrow \\ \text{time evolution}}}{B(t)}$$

$$\Rightarrow A(x) B''(t) = B(t) A''(x)$$

$$\Rightarrow \frac{B''(t)}{B(t)} = \frac{A''(x)}{A(x)} = \text{constant} = -\omega^2$$

We get two equations: $B''(t) = -\omega^2 B(t)$.

$$A''(x) = -\omega^2 A(x)$$

$$\Rightarrow B(t) = \alpha \cos(\omega t) + \beta \sin(\omega t)$$

$$A(x) = a \cos(\omega x) + b \sin(\omega x)$$

Boundary condition: $A(0) = A(\pi) = 0 \Rightarrow a = 0$

$$\Rightarrow \psi_\omega(t, x) = b \sin(\omega x) \left[\underbrace{\alpha_\omega \cos(\omega t) + \beta_\omega \sin(\omega t)}_{A_\omega(t)} \right]$$

By superposition, formally.

$$\psi(t, x) = \sum_{\omega > 0}^{\infty} A_\omega(t) \sin(\omega x). \text{ solves (eq)}$$

where $A_\omega(t) = \underline{\alpha_\omega} \cos(\omega t) + \underline{\beta_\omega} \sin(\omega t)$

By initial conditions

Consider initial conditions: $\psi(0, x) = f(x)$

$$\partial_t \psi(0, x) = g(x)$$

we require that $\sum_{\omega > 0}^{\infty} \alpha_\omega \sin(\omega x) = f(x)$

$$\sum_{\omega > 0}^{\infty} \omega \beta_\omega \sin(\omega x) = g(x)$$

This raises the question:

Given f on $[0, \pi]$, $f(0) = f(\pi) = 0$. Can we find

α_n such that $f(x) = \sum_{n=0}^{\infty} \alpha_n \sin(n x)$?

Theorem 1: Let $f \in L^2(\mathbb{T})$. Then $\lim_{N \rightarrow \infty} \|S_N f - f\|_{L^2(\mathbb{T})} = 0$.

Remark: $S_N f(x) = \sum_{|k| \leq N} \widehat{f}(k) e^{2\pi i k x}$, $S_N g(x) = \sum_{|k| \leq N} \widehat{g}(k) e^{2\pi i k x}$

$$\begin{aligned} & \int_0^1 S_N f(x) \overline{S_N g(x)} dx \\ &= \int_0^1 \sum_{|k| \leq N} \widehat{f}(k) e^{2\pi i k x} \sum_{|m| \leq N} \overline{\widehat{g}(m)} e^{-2\pi i m x} dx \\ &= \sum_{|k|, |m| \leq N} \widehat{f}(k) \overline{\widehat{g}(m)} \underbrace{\int_0^1 e^{2\pi i (k-m)x} dx}_{\delta_{km}} \end{aligned}$$

$$= \sum_{|k| \leq N} \widehat{f}(k) \overline{\widehat{g}(k)}$$

In particular, $\|S_N f\|_{L^2(\mathbb{T})}^2 = \sum_{|k| \leq N} |\widehat{f}(k)|^2$

pf: 1° First we observe that $e_k(x) = e^{2\pi i k x}$ form an orthonormal basis of $L^2(\mathbb{T})$

2°. $S_N f = \sum_{|k| \leq N} (f, e_k) e_k$

therefore, $(f - S_N f, e_j) = 0 \quad \forall |j| \leq N$

Then the orthonormal property of the family $\{e_k\}$ implies that $f - S_N f$ is orthogonal to e_j for all $|j| \leq N$. i.e.

$$(f - S_N f) \perp \text{span} \{e_j; |j| \leq N\}$$

In particular,

$$\begin{aligned} \|f\|_{L^2}^2 &= \|f - S_N f\|_{L^2}^2 + \|S_N f\|_{L^2}^2 \\ &\geq \|S_N f\|_{L^2}^2 \\ &= \sum_{|k| \leq N} |(f, e_k)|^2 \end{aligned}$$

$$\Rightarrow S_N f \xrightarrow{N \rightarrow \infty} \sum_{k \in \mathbb{Z}} (f, e_k) e_k \quad \text{in } L^2$$

$$\Rightarrow \lim_{N \rightarrow \infty} \|S_N f - f\|_{L^2} = 0$$

□

Pf of " $\widehat{f * g}(\zeta) = \widehat{f} \widehat{g}$ "

$$\begin{aligned}\widehat{f * g}(\zeta) &= \frac{1}{2\pi} \int_{-2}^2 (f * g)(x) e^{-ix\zeta} dx \\ &= \frac{1}{2\pi} \int_{-2}^2 \frac{1}{2\pi} \int_{-2}^2 f(y) g(x-y) dy e^{-ix\zeta} dx \\ &= \frac{1}{2\pi} \int_{-2}^2 f(y) e^{-i\zeta y} \left(\frac{1}{2\pi} \int_{-2}^2 g(x-y) e^{-i\zeta(x-y)} dx \right) dy \\ &= \frac{1}{2\pi} \int_{-2}^2 f(y) e^{-i\zeta y} \left(\frac{1}{2\pi} \int_{-2}^2 g(x) e^{-i\zeta x} dx \right) dy \\ &= \widehat{f}(\zeta) \widehat{g}(\zeta)\end{aligned}$$